



SM parameters from Lattice QCD

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HPQCD collaboration

EFT & LGT
Munich, May 2016

Quark masses and strong coupling are fundamental parameters of the SM but cannot be directly determined from experiment.

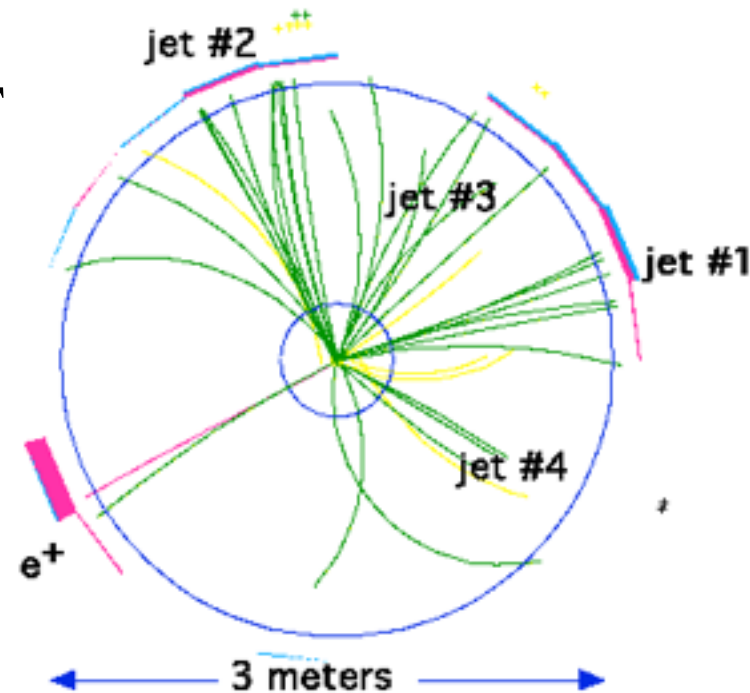
Well-defined m_q/α_s are scheme and scale-dependent.

Convention to use \overline{MS}

Compare results from multiple approaches for strong test of QCD.

Masses are input to theoretical expressions for SM cross-sections e.g. $H \rightarrow c\bar{c}$

CDF



	Higgs X-Section WG	PDG	lattice	Karlsruhe (e ⁺ e ⁻)	world non-lattice
$\delta \alpha_s$	0.002	0.0007	0.0007		0.0012
δm_c (GeV)	0.03	0.025	0.006	0.013	
δm_b (GeV)	0.06	0.03	0.023	0.016	

P. Mackenzie, Snowmass 2013

Conversion of lattice quark masses to \overline{MS} scheme

- Direct methods: Determine $m_{q,latt}$ in lattice QCD.

$$m_{\overline{MS}}(\mu) = Z_m(\mu a) m_{latt}$$

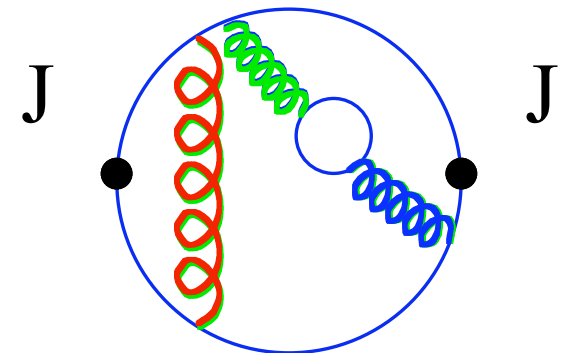
Calculate Z in lattice QCD pert. th. or use ‘nonpert’ RI-MOM lattice matching.

Error dominated by that of Z and continuum extrapolation.

Note: Z cancels in mass ratios.

- Indirect methods: (after tuning m_{latt}) match a quantity calculated in lattice QCD to continuum pert. th. in terms of \overline{MS} quark mass

e.g. Current-current correlators for heavy quarks known through α_s^3 .



Issues with handling ‘heavy’ quarks on the lattice:

$$L_q = \bar{\psi}(\not{D} + m)\psi \rightarrow \bar{\psi}(\gamma \cdot \Delta + ma)\psi$$

Δ is a finite difference on the lattice - leads to discretisation errors. What sets the scale for these?

For light hadrons the scale is Λ_{QCD} = few hundred MeV

For heavy hadrons the scale can be m_Q

$$E(a) = E(a = 0) \times (1 + A(m_Q a)^2 + B(m_Q a)^3 + \dots)$$

$$m_c a \approx 0.4, m_b a \approx 2 \quad \text{for } a \approx 0.1\text{fm}$$

➡ need good discretisation of Dirac equation and multiple values of a for accurate continuum extrapolation.

Highly Improved Staggered Quarks (HISQ) formalism has errors improved to $\alpha_s(am)^2, (am)^4$

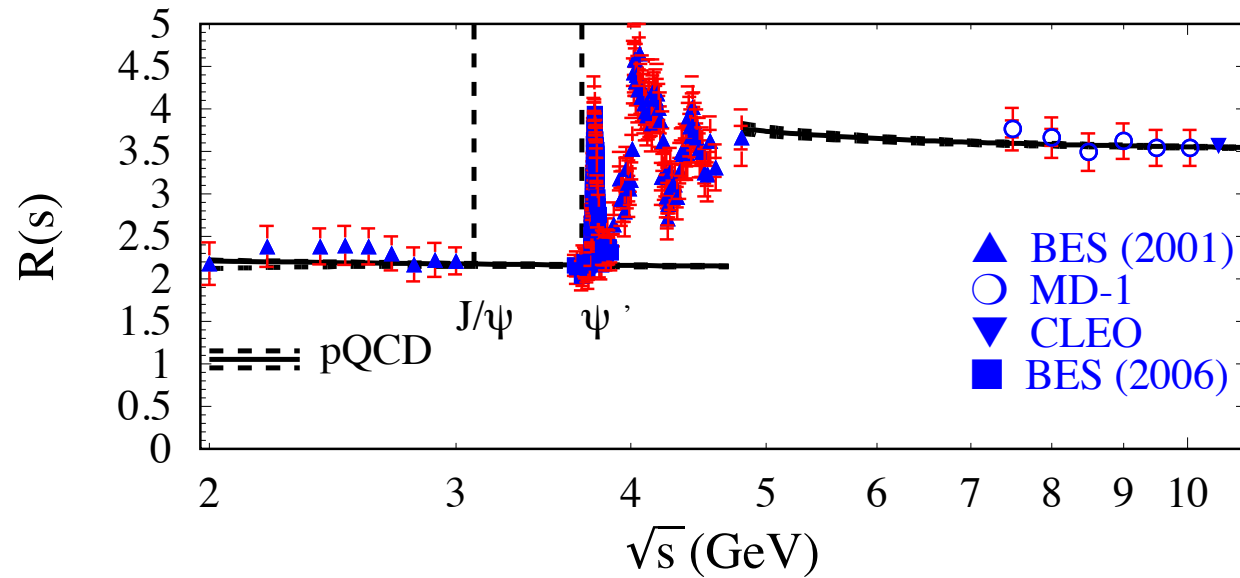
Follana et al, HPQCD,
hep-lat/0610092

Current-current correlator method for m_c

Continuum: extract charm piece of:

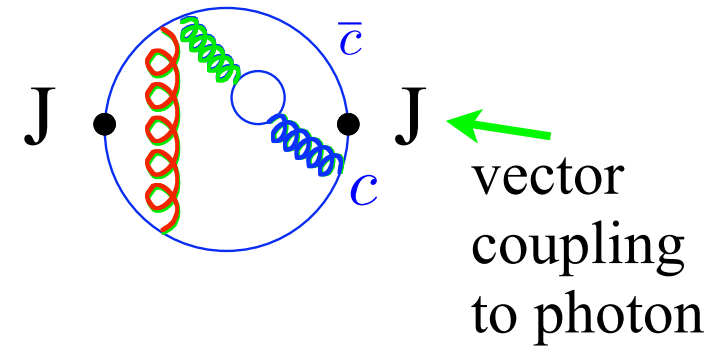
e.g. Kuhn et al,
hep-ph/0702103

$R_{e^+e^-}(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{4\pi\alpha^2/(3s)}$ from experiment, then



$$\mathcal{M}_k \equiv \int \frac{ds}{s^{k+1}} R_{e^+e^-}(s)$$

$$= \frac{12\pi^2}{n!} \left(\frac{d}{dq^2} \right)^k \Pi_c(q^2) \Big|_{q^2=0}$$



$$\Pi_c(q^2) = \frac{3}{16\pi^2} e_c^2 \sum_{k \geq 0} C_k^V \left(\frac{q^2}{4(m_c(\mu))^2} \right)^k$$

C_k a power series in $\alpha_s(\mu)$, known through α_s^3 for first few values of k

Use $k=1$: $m_c(m_c) = 1.279(13)\text{GeV}$

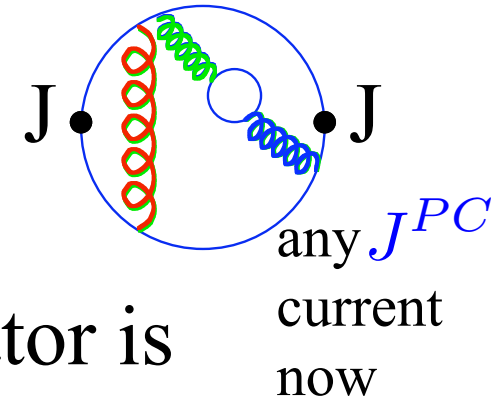
errors: expt + α_s

Chetyrkin et al,
0907.2110

Current-current correlator method for lattice m_c

HPQCD + Chetyrkin et al, 0805.2999, C. Mcneile et al, HPQCD,1004.4285

• Substitute time-moment of lattice charmonium correlator for experiment. In principle can use any current J now.



• For HISQ quarks pseudoscalar η_c correlator is most accurate. J is absolutely normalised.

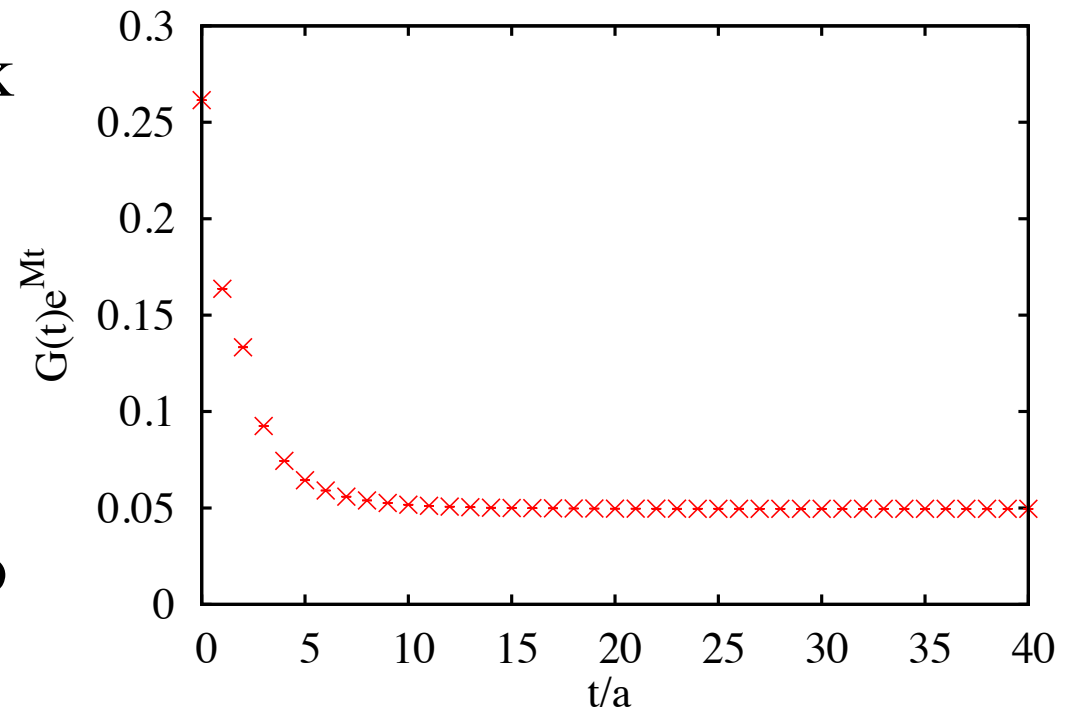
step 1: calculate η_c correlators by combining lattice charm quark propagators

step 2: large time - fit to exponential, gives η_c mass

step 3: tune lattice quark mass so η_c mass correct.

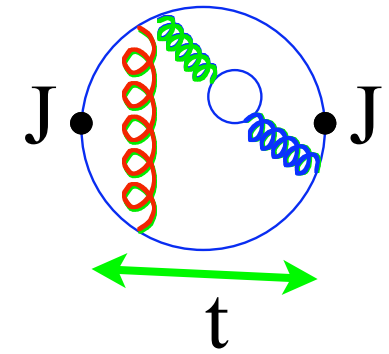
step 4: calculate time moments to compare to QCD pert. theory.

Emphasises short-time contribns.



Correlator time-moments:

$$G(t) = a^6 \sum_{\vec{x}} (am_c)^2 \langle 0 | j_5(\vec{x}, t) j_5(0, 0) | 0 \rangle$$



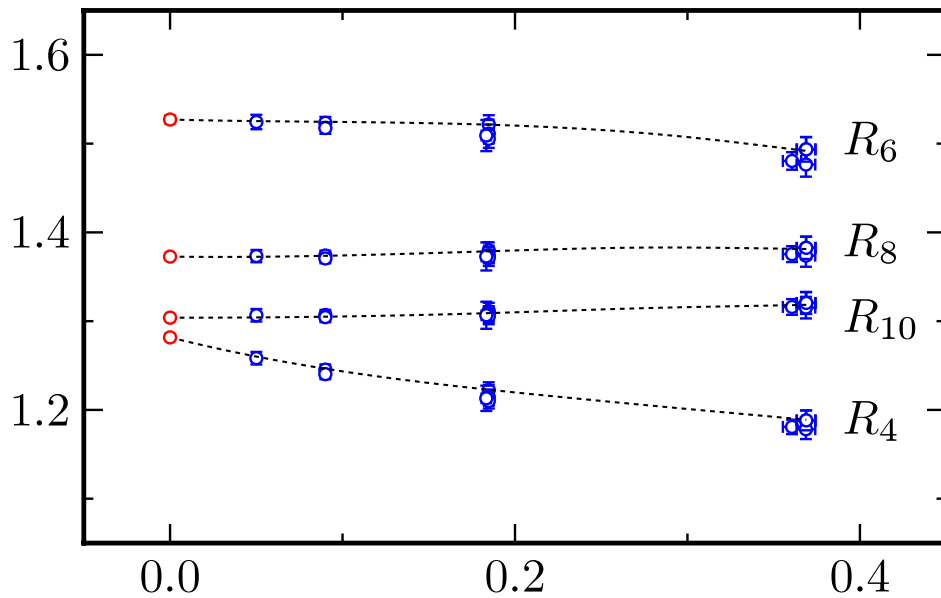
$$G_n = \sum_t (t/a)^n G(t)$$

$$R_{n,latt} = G_4 / G_4^{(0)} \quad n = 4$$

ratio to results with no gluon field improves disc. errors

$$= \frac{am_{\eta_c}}{2am_c} (G_n / G_n^{(0)})^{1/(n-4)} \quad n = 6, 8, 10 \dots$$

(match $k = 2, 3, 4 \dots$)



extrapolate to $a=0$ and compare to contnm pert. th.

$$R_{n,cont} = \frac{m_{\eta_c}}{2m_c(\mu)} \frac{C_k^P}{C_k^{P,0}} \quad n = 2k + 2$$

$$\frac{C_k^P}{C_K^{P,0}} = 1 + \sum c_i \alpha_s^i(\mu)$$

$n_f = 2 + 1$

$a^2 \text{ (GeV}^{-2}\text{)}$

Fit first 4 moments simultaneously, gives

$$\frac{m_{\eta_c}}{2m_c(\mu)} \quad \text{AND} \quad \alpha_s(\mu)$$

$\mu = 3m_c$

Result:

$$m_c(m_c) = 1.273(6) \text{ GeV}$$

error dominated by unknown higher orders in pert. th.

C. McNeile et al, HPQCD, 1004.4285

Further check:

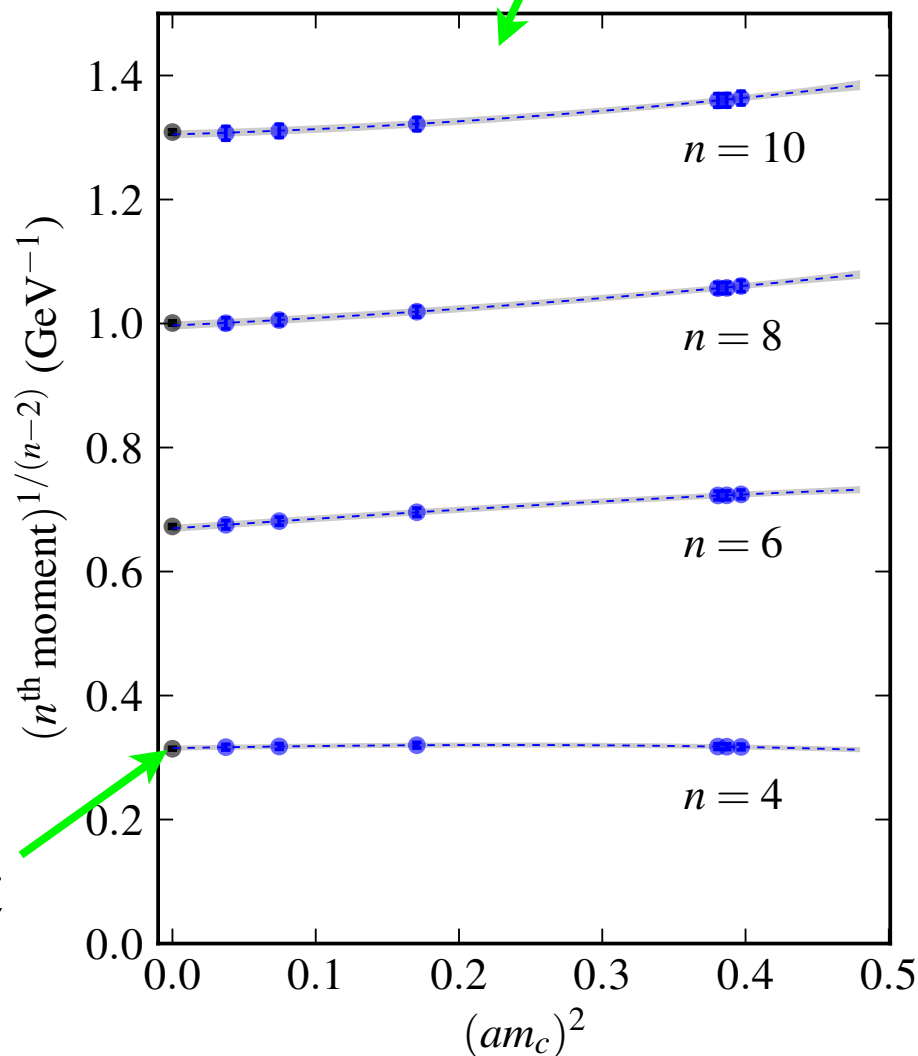
compare vector moments (after normalising current) to those extracted from $R_{e^+e^-}$

Agreement is a 1% test of (lattice) QCD. Gives:

$$a_\mu^c = 14.4(4) \times 10^{-10}$$

HPQCD, 1208.2855, 1403.1778

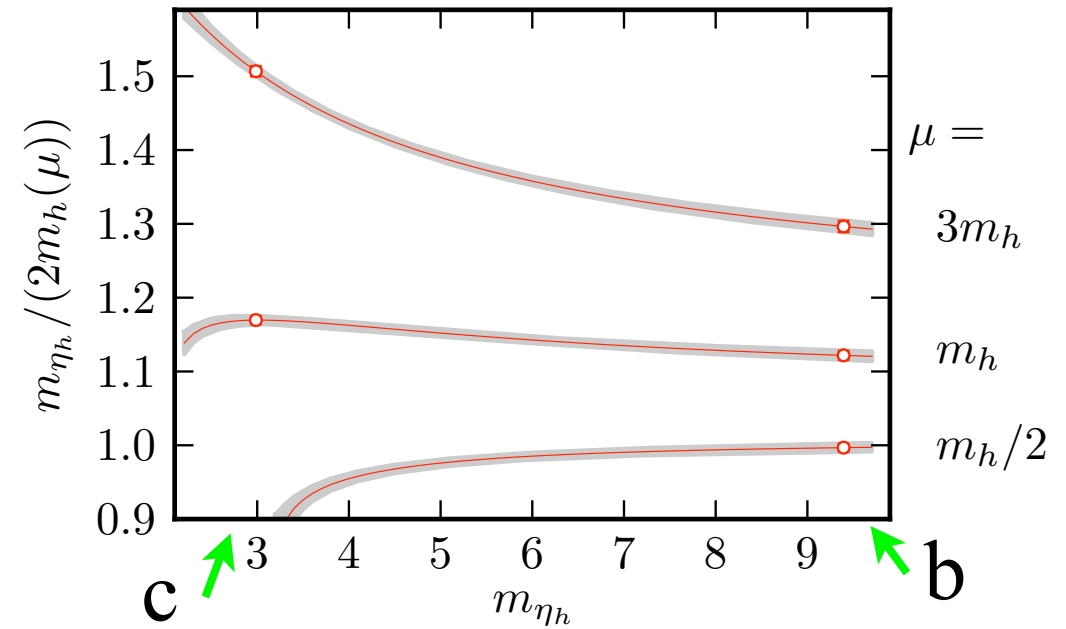
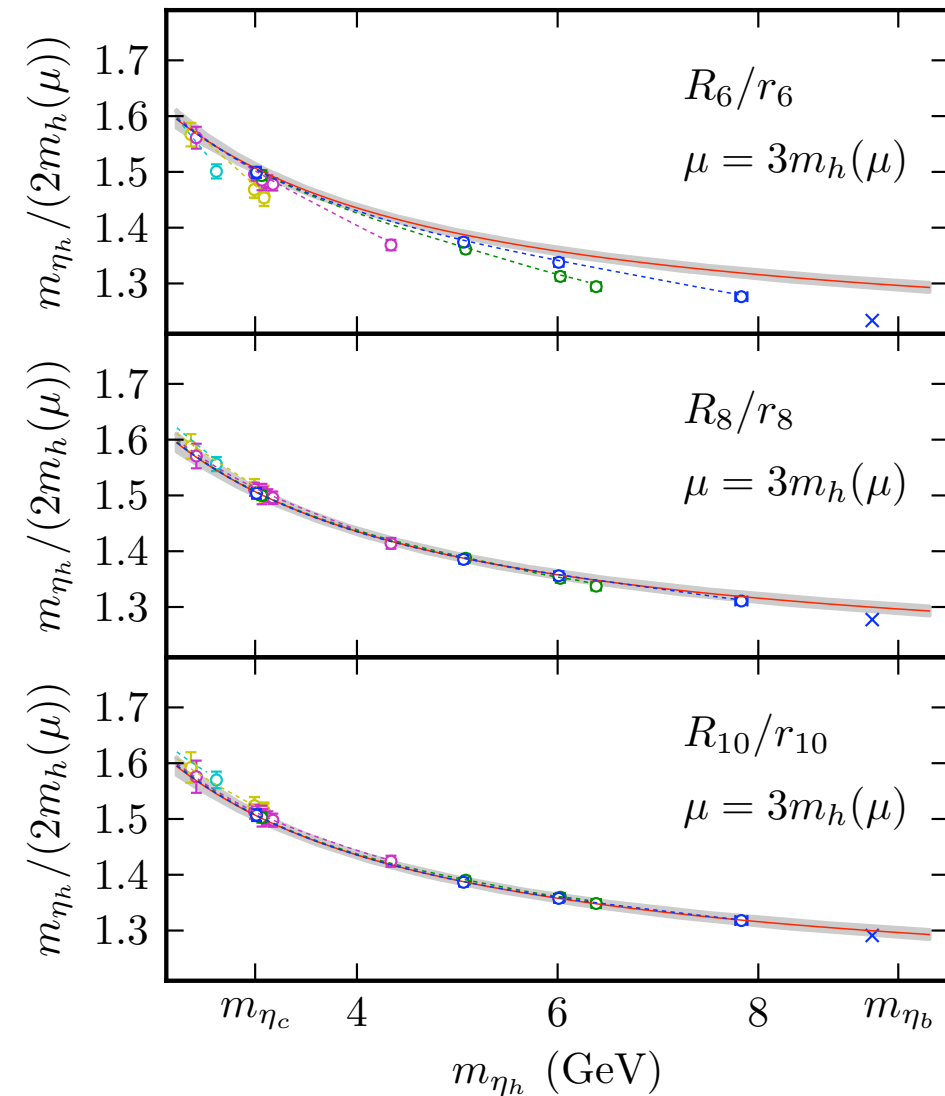
lattice and expt errors similar size



expt

see also ETMC, 1111.5252

- Repeat calcln for $m_q \geq m_c$ inc. ultrafine lattices



Can determine m_h/m_{η_h} for heavy quarks - extrapolate (slightly) to b.

$$\overline{m}_b^{n_f=5}(\overline{m}_b) = 4.164(23)\text{GeV}$$

key error is now extrapoln in a

Agrees well with contnm results using R_{e+e-}

Update and improved method

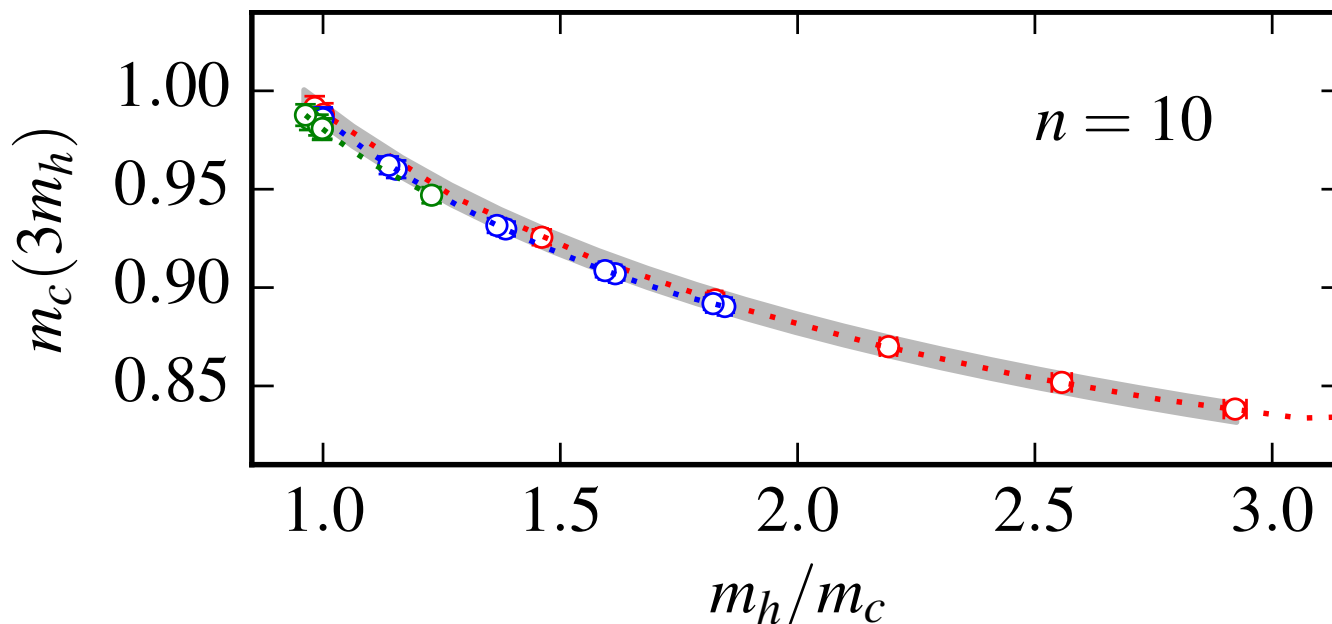
HPQCD, 1408.4169

Use improved $n_f = 2+1+1$ gluon field configs, more accurate lattice spacing determination etc etc.

Determine m_c at higher scales by using multiple m_h

$$\tilde{R}_n = \frac{1}{m_c} \left(G_n / G_n^{(0)} \right)^{1/(n-4)} \rightarrow \frac{1}{m_c(\mu)} \frac{C_k^P}{C_k^{P.0}}$$

\swarrow tuned m_c \swarrow G_n at m_h \swarrow $\mu = 3m_h$

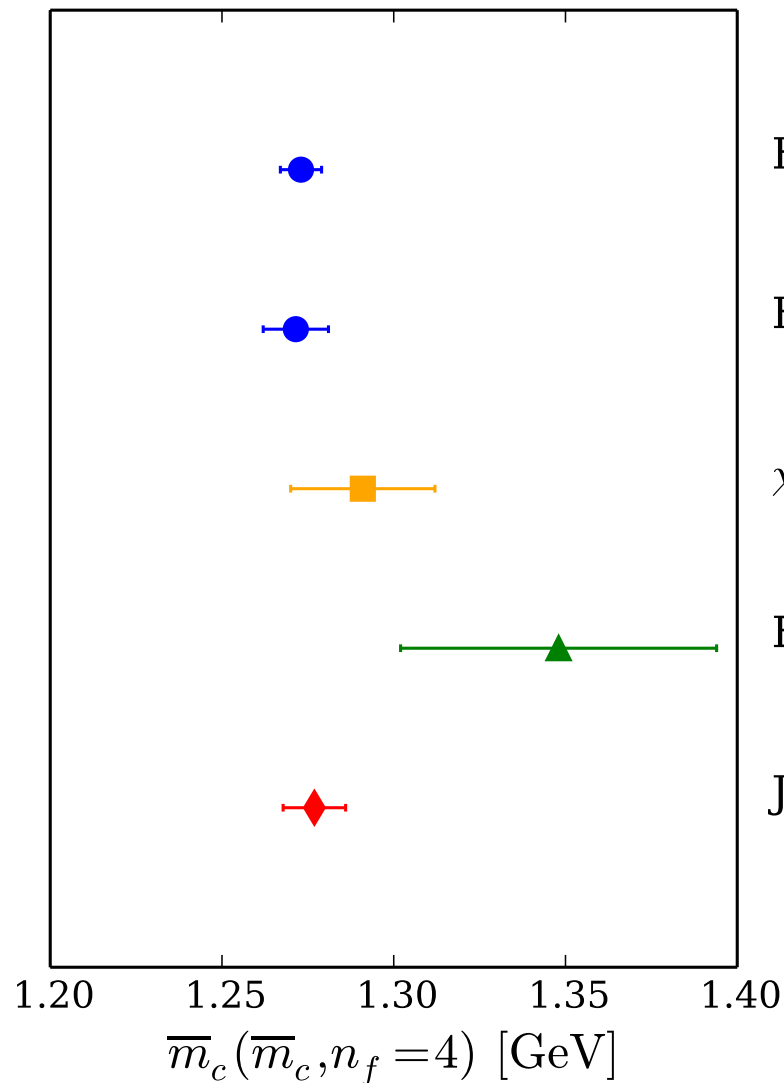


$$\frac{m_{0h}}{m_{0c}} = \frac{m_h(\mu)}{m_c(\mu)}$$

$$m_c(m_c) = 1.2715(95) \text{ GeV}$$

m_c summary

Good consistency between lattice methods and actions



HPQCD HISQ $n_f=3$ [1004.4285] JJ

HPQCD HISQ $n_f=4$ [1408.4169] JJ

χ QCD $n_f=3$ [1410.3343] RI-MOM

ETMC $n_f=4$ [1403.4504] RI-MOM

JLQCD $n_f=3$ [1511.09163] JJ

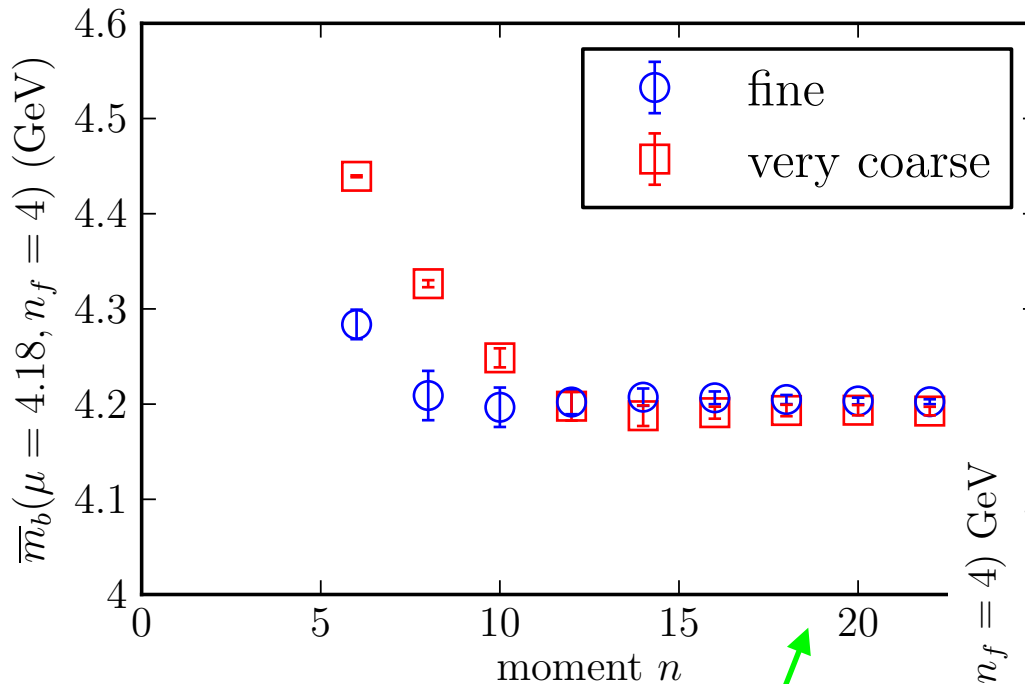
Alternative determinations of m_b

HPQCD, 1408.4768

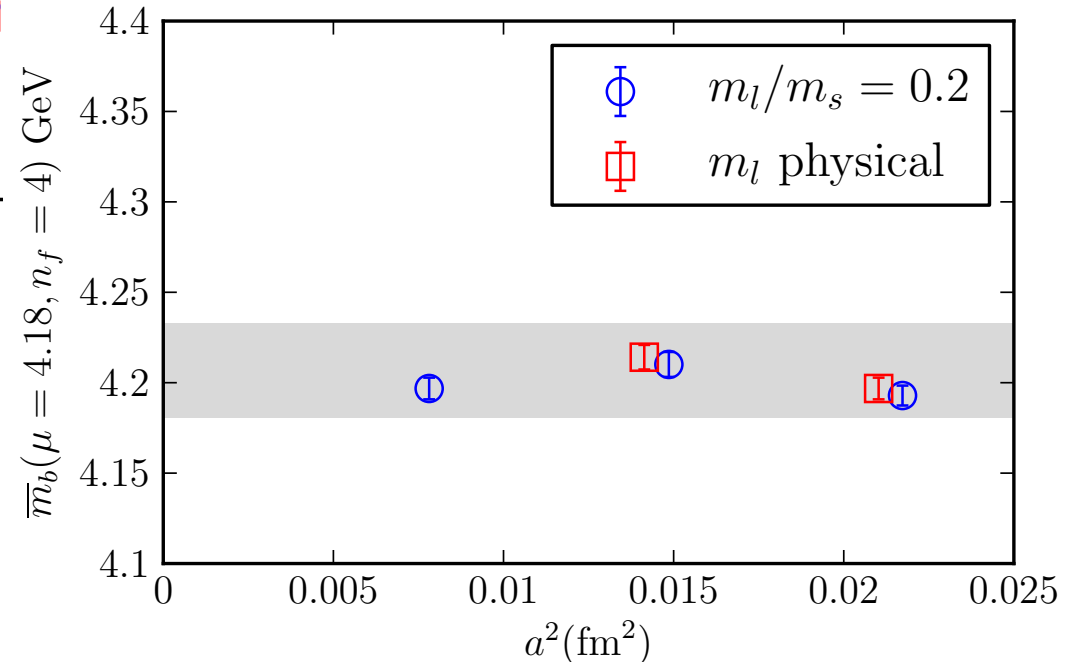
Current-current correlator method using vector bottomonium correlators calculated with improved NRQCD b quarks

$n_f = 2+1+1$ HISQ sea quarks

$$\bar{m}_b(\bar{m}_b, n_f = 5) = 4.196(25)\text{GeV}$$

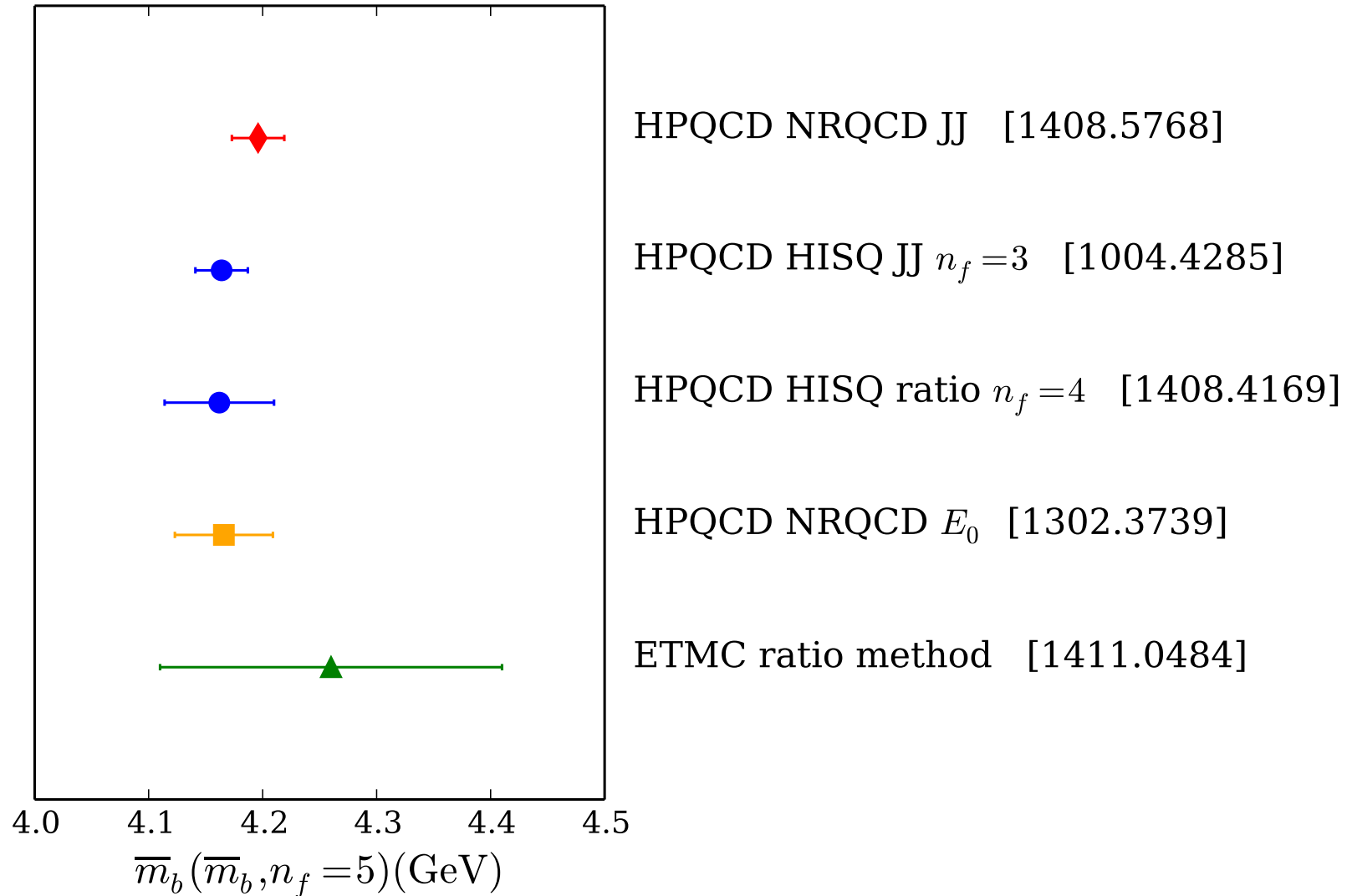


Larger moment numbers more nonrelativistic - use 18



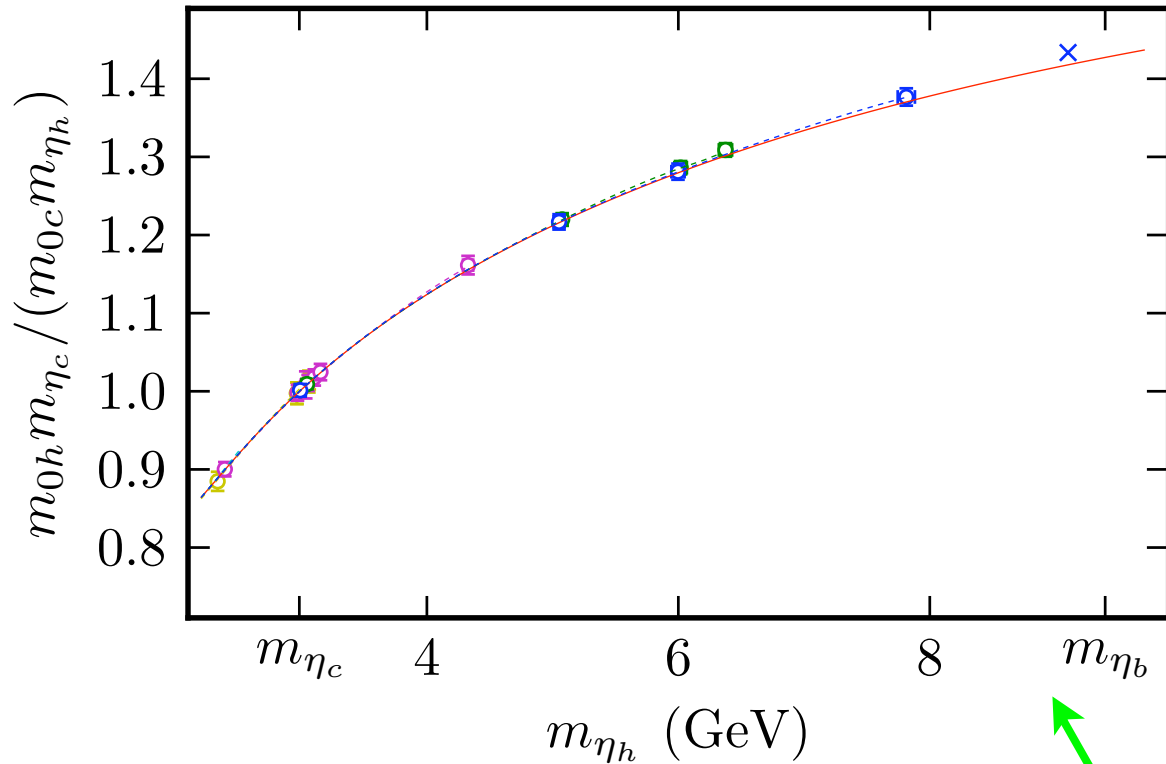
m_b summary

Again, good consistency between different lattice methods



m_b/m_c from lattice QCD

HPQCD, 1004.4285



$$\left(\frac{m_{q1,latt}}{m_{q2,latt}} \right)_{a=0} = \frac{m_{q1,\overline{MS}}(\mu)}{m_{q2,\overline{MS}}(\mu)}$$

completely nonperturbative determination of ratio gives:

$$\frac{m_b}{m_c} = 4.49(4)$$

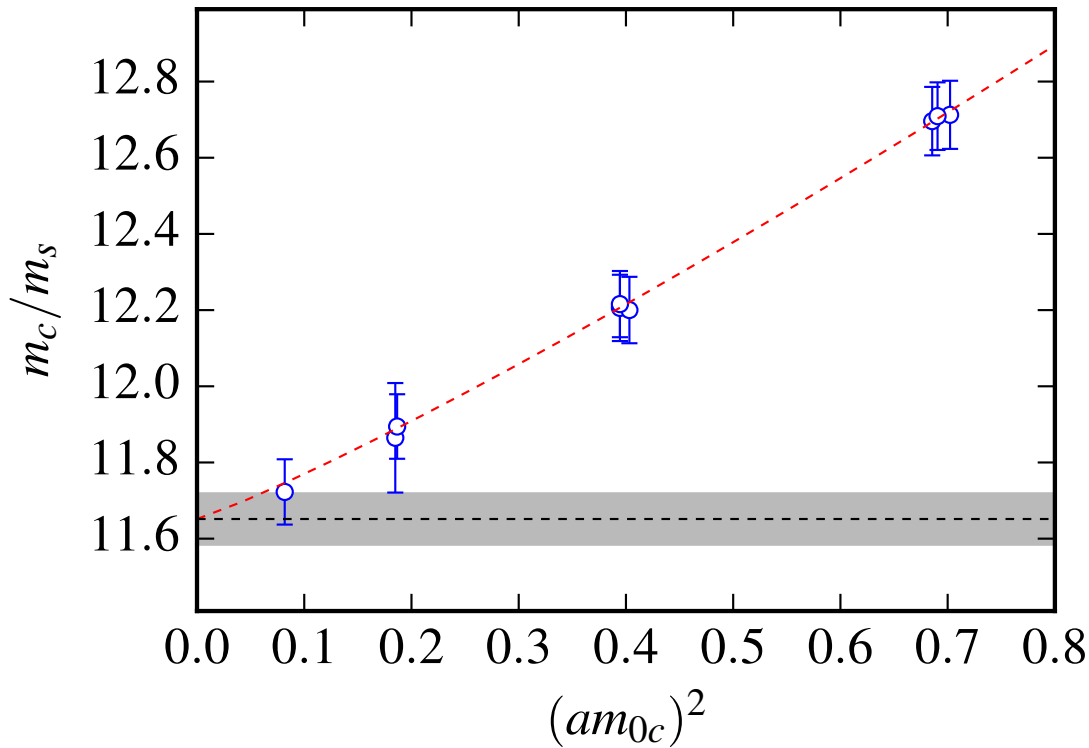
Agrees with that from current-current correlator method - test of pert. th. . Also tested $n_f=2+1+1$

HPQCD,
1408.4169

$$m_c/m_s$$

Mass ratio can be obtained directly from lattice QCD if same quark formalism is used for both quarks.

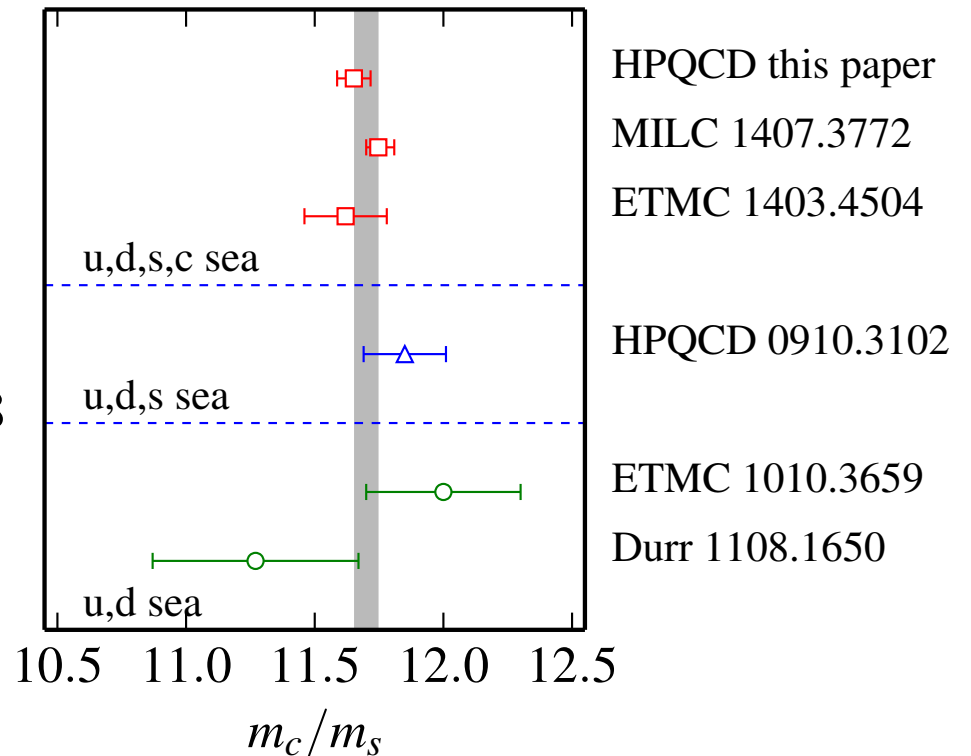
Not possible with any other method ...



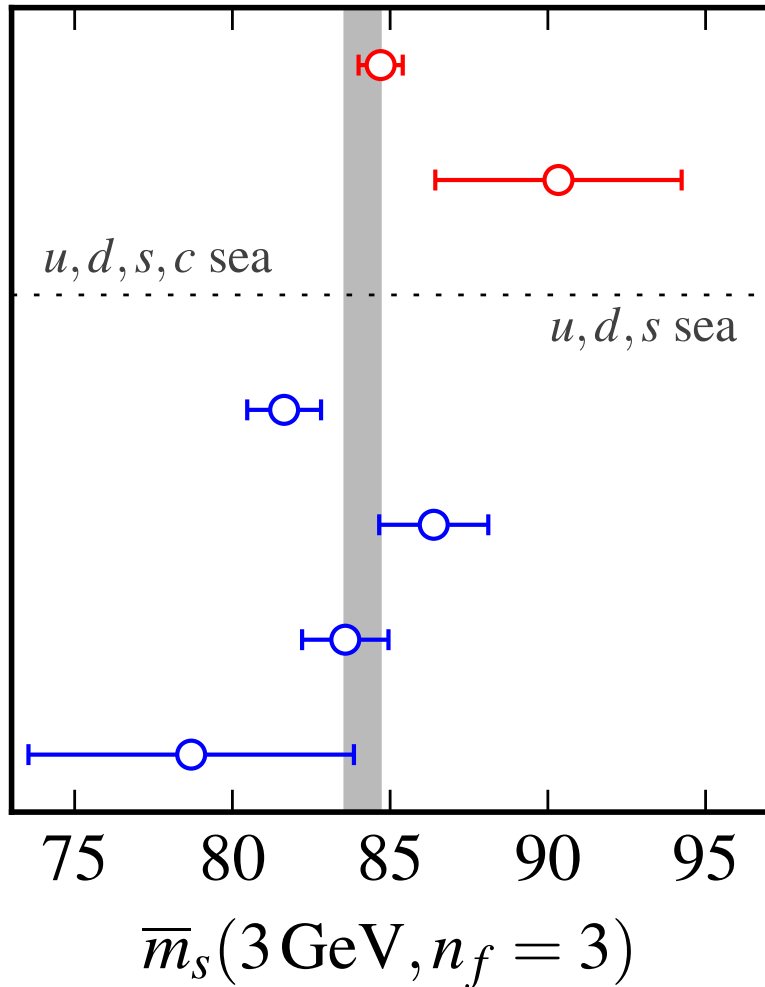
$$\frac{m_c}{m_s} = 11.652(65)$$

HPQCD, 1408.4169

Good consistency
between different lattice
actions



Combining m_c and m_c/m_s leads to 1% accuracy in m_s



HPQCD 1408.4169

ETMC 1403.4504 RI-MOM

RBC/UKQCD 1411.7017 RI-SMOM

Durr et al 1011.2403 RI-MOM

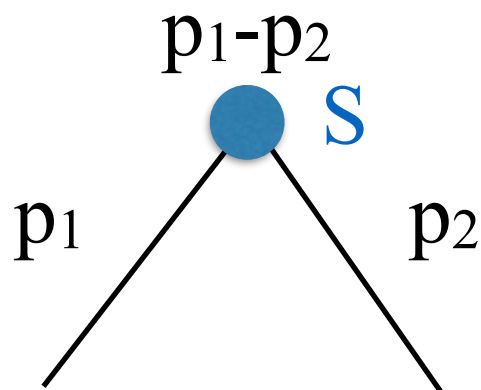
HPQCD 0910.3102

HPQCD (pert) 0511160

$$\bar{m}_s(3 \text{ GeV}, n_f = 3) = 84.1(5) \text{ MeV}$$

New HPQCD results for m using RI-SMOM scheme

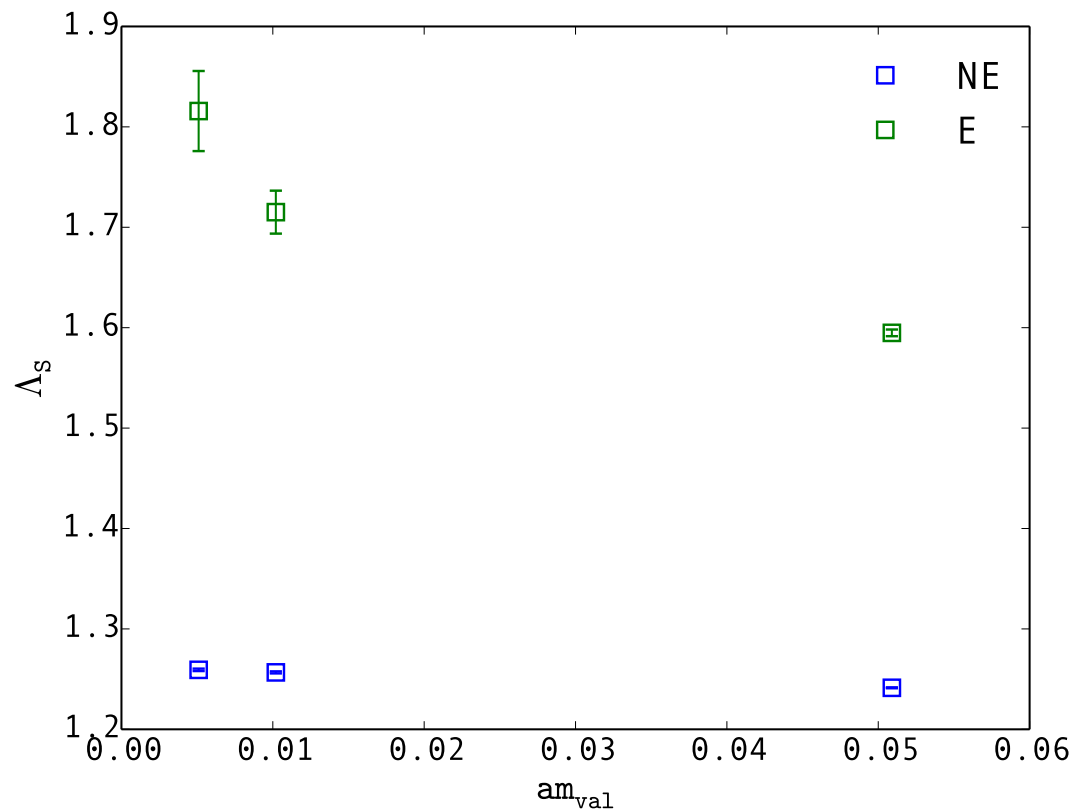
Direct test of mass determination using different method to current-current correlators but the same formalism



Use RI-MOM method, fixing Landau gauge Green's function of quark bilinear to tree-level value. Match to \overline{MS} perturbatively.

$$Z_m = Z_S^{-1}$$

Lytle, HPQCD, 1511.06547



‘Non-exceptional’ kinematics:

$$p_1^2 = p_2^2 = (p_1 - p_2)^2$$

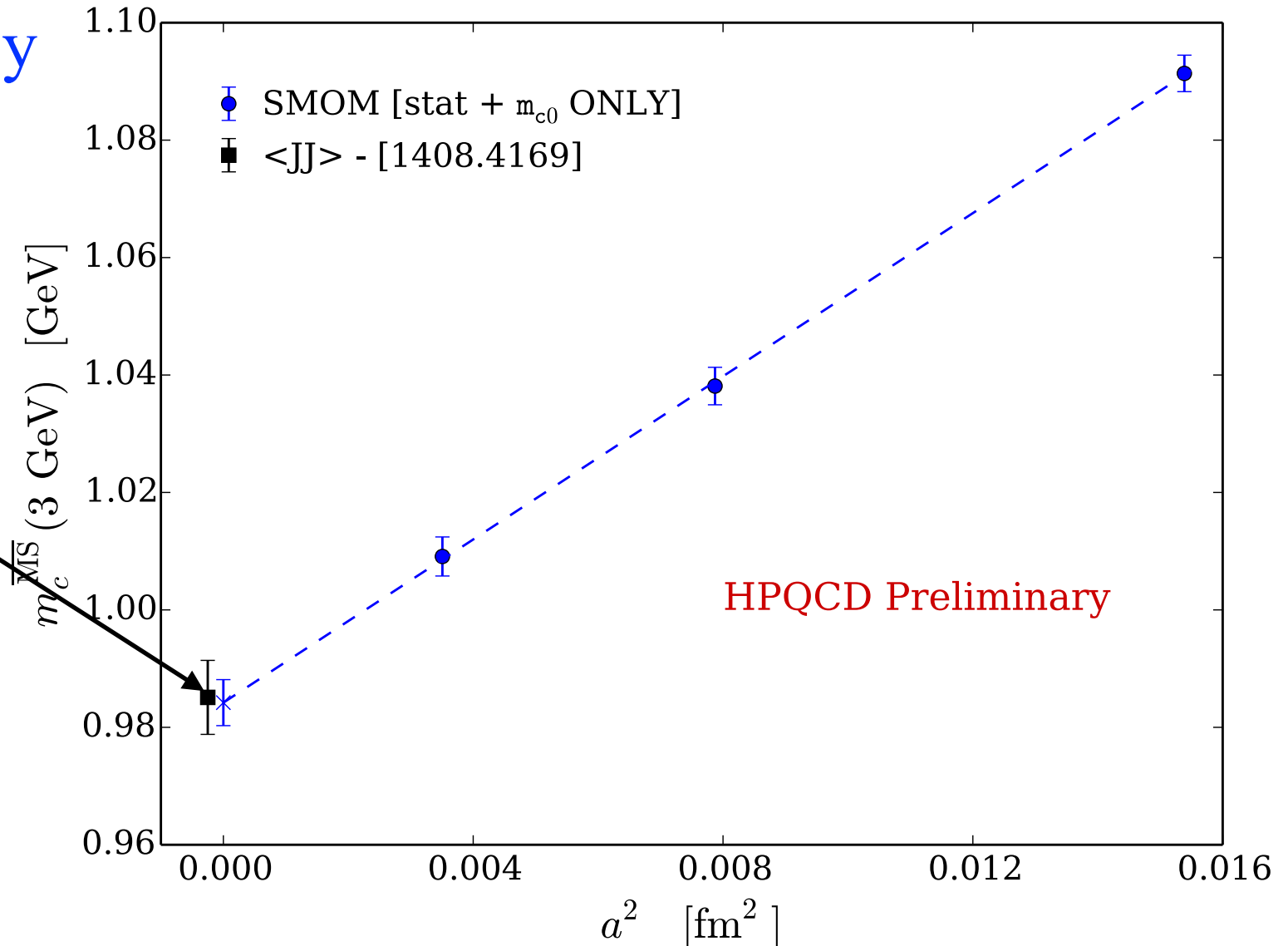
much improved results Sturm et al, 0901.2599

Infrared sensitivity much reduced and perturbative matching factor (known through α_s^2) very close to 1.

Multiply tuned bare masses by Z_m :

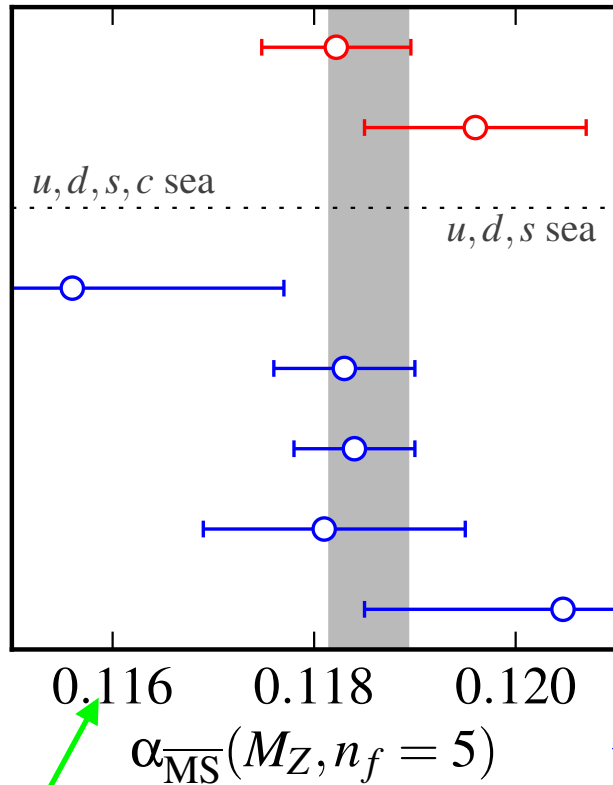
Good agreement with results from current-current correlators

still to do: error budget



α_s summary

α_s is also well determined in current-current correlator method (particularly lowest moment)



HPQCD-*jj* 1408.4169

ETMC 1310.3763 *ghost-gluon*

static potential

Basavov et al 1205.6155

HPQCD-*jj* 1004.4285

HPQCD-*W_{nm}* 1004.4285

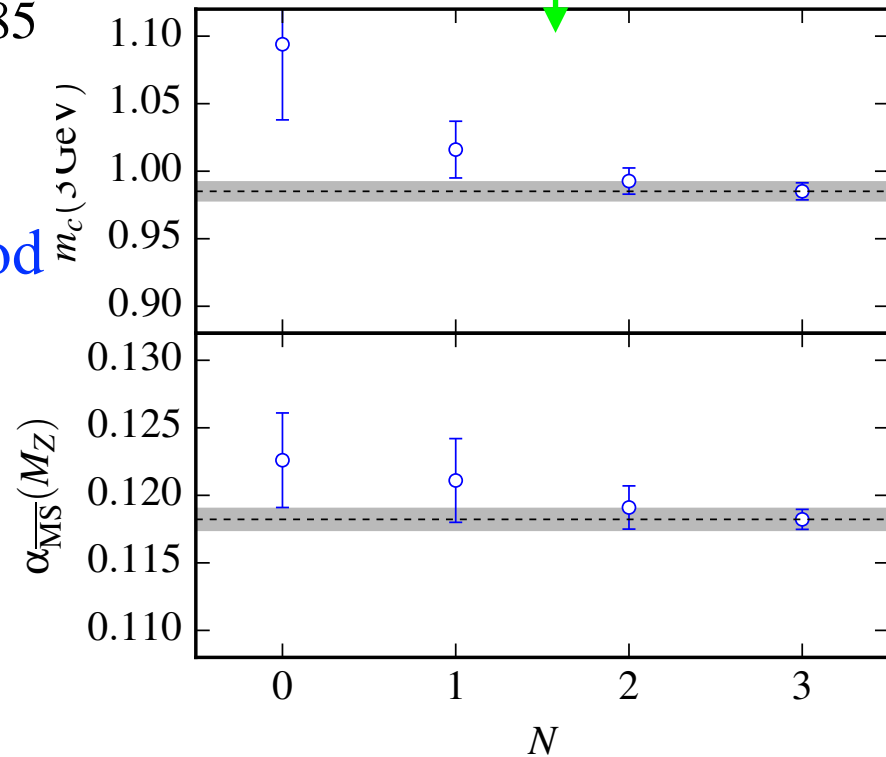
JLQCD 1002.0371

PACS-CS 0906.3906

SF method

vacuum polarisation

Effect of high-order pert. theory in JJ method



Good consistency from very different lattice methods

$$\alpha_{\overline{MS}}(M_Z, n_f = 5) = 0.1185(4)$$

Conclusions

$m_c(m_c)$ is determined to 1% and
 $m_b(m_b)$ to 0.5% from continuum and lattice methods.
 $\alpha_S(M_Z)$ to 0.5% from lattice - multiple methods

1% accurate m_c/m_s ratio allows 1% in m_s also, along
with RI-MOM methods

Tests of perturbation theory from completely non-
perturbative mass ratios and JJ/RI-MOM comparison

Future improvements from higher order pert th. (?possible)
and finer lattices to push up mu values.

Error budget for HISQ current-current method

TABLE IV. Error budget [31] for the c mass, QCD coupling, and the ratios of quark masses m_c/m_s and m_b/m_c from the $n_f = 4$ simulations described in this paper. Each uncertainty is given as a percentage of the final value. The different uncertainties are added in quadrature to give the total uncertainty. Only sources of uncertainty larger than 0.05% have been listed.

HPQCD,
1408.4169

	$m_c(3)$	$\alpha_{\overline{\text{MS}}}(M_Z)$	m_c/m_s	m_b/m_c
Perturbation theory	0.3	0.5	0.0	0.0
Statistical errors	0.2	0.2	0.3	0.3
$a^2 \rightarrow 0$	0.3	0.3	0.0	1.0
$\delta m_{uds}^{\text{sea}} \rightarrow 0$	0.2	0.1	0.0	0.0
$\delta m_c^{\text{sea}} \rightarrow 0$	0.3	0.1	0.0	0.0
$m_h \neq m_c$ (Eq. (15))	0.1	0.1	0.0	0.0
Uncertainty in $w_0, w_0/a$	0.2	0.0	0.1	0.4
α_0 prior	0.0	0.1	0.0	0.0
Uncertainty in m_{η_s}	0.0	0.0	0.4	0.0
$m_h/m_c \rightarrow m_b/m_c$	0.0	0.0	0.0	0.4
δm_{η_c} : electromag., annih.	0.1	0.0	0.1	0.1
δm_{η_b} : electromag., annih.	0.0	0.0	0.0	0.1
Total:	0.64%	0.63%	0.55%	1.20%