Heavy Flavor Physics: BSM phenomenology



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Outline

- Motivation and introduction
 for BSM focus (mostly) on loop processes
- LQCD results for
 - \bigstar semileptonic *B* meson form factors
 - \Rightarrow neutral *B* meson mixing matrix elements
 - ☆ summary of recent progress
- Phenomenology
 - ☆ SM pre/post-dictions
 - ☆ CKM unitarity & BSM implications
 - ☆ Lepton Flavor Universality
- Summary

Motivation

example: $B^+ \to K^+ \ell^+ \ell^-$



Experiment vs. SM theory:



Semileptonic *B*-meson decay to light hadrons



★ shape for semileptonic *B* decays:

use z-expansion for model-independent parameterization of q^2 dependence \star calculate the complete set of form factors, $f_+(q^2)$, $f_0(q^2)$ and $f_T(q^2)$ with LQCD.

Rare semileptonic B decay



Parameterize the amplitude in terms of the three form factors $f_{+,0,T}(q^2)$:

$$A(B \to P \,\ell \ell) \sim C_7^{\text{eff}} f_T + (C_9^{\text{eff}} + C_{10}) f_+ + \text{nonfactorizable terms}$$

more on these later



Form factors for $B \to K \ell \ell$



HPQCD (arXiv:1306.0434, 1306.2384, PRL 2013)

FNAL/MILC (arXiv:1509.06235, PRD 2016)

 \Rightarrow Two LQCD calculations (on overlapping ensemble sets, different valence actions): HPQCD (NRQCD *b* + HISQ), FNAL/MILC (Fermilab *b* + asqtad)

- ☆ consistent results for all three form factors
- * also consistent with LCSR (Khodjamarian et al, arXiv:1006.4945, JHEP 2010)
- ☆ Note: First LQCD calculation of $\Lambda_b \rightarrow \Lambda \ell^+ \ell^-$ form factors (10 total) (Detmold and Meinel, arXiv:1602.01399, PRD 2016)



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Form factors for $B \to \pi \, \ell \ell$



- \Rightarrow Two independent calculations (RBC, FNAL/MILC): consistent results for f_0, f_+
- \Rightarrow Shape of f_+ agrees with experiment and uncertainties are commensurate
- ★ Fit lattice form factors together with experimental data to determine $|V_{ub}|$ and improved form factors (f_{+}, f_0)
- \Rightarrow **First** calculation of $f_{\rm T}$ (FNAL/MILC)



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Form factors for $B \to D \,\ell \nu, \ \ell = e, \mu, \tau$



HPQCD (arXiv:1505.03925, PRD 2015) FNAL/MILC (arXiv:1503.07237, PRD 2015)

- ☆ Two LQCD calculations (FNAL/MILC, HPQCD)
- Combined fit to LQCD form factors and BaBar data.
- LQCD form factor
 uncertainties (~1.2%)
 smaller than experiment.

 \Rightarrow Form factors can be used to calculate the CKM free ratio:

$$R(D) \equiv \frac{\mathcal{B}(B \to D\tau\nu_{\tau})}{\mathcal{B}(B \to D\ell\nu)}$$

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In general :	SM:	BSM:
$\mathcal{H}_{\text{eff}} = \sum_{i=1}^{5} c_i(\mu) \mathcal{O}_i(\mu)$	$\mathcal{O}_1 = (\bar{b}^{\alpha} \gamma_{\mu} L q^{\alpha}) \; (\bar{b}^{\beta} \gamma_{\mu} L q^{\beta})$	$\mathcal{O}_4 = (\bar{b}^{\alpha} L q^{\alpha}) \; (\bar{b}^{\beta} R q^{\beta})$
	$\mathcal{O}_2 = (\bar{b}^{\alpha} L q^{\alpha}) \; (\bar{b}^{\beta} L q^{\beta})$	$\mathcal{O}_5 = (\bar{b}^{\alpha} L q^{\beta}) \; (\bar{b}^{\beta} R q^{\alpha})$
	$\mathcal{O}_3 = (\bar{b}^{\alpha} L q^{\beta}) \; (\bar{b}^{\beta} L q^{\alpha})$	

Recent and ongoing LQCD calculations of K, D, and B mixing quantities now include results for hadronic matrix elements of all five operators.



☆ new LQCD calculation by FNAL/MILC (Fermilab *b* + asqtad)
 ☆ significant reduction of errors, especially for ξ
 ☆ first three flavor LQCD result for all five matrix elements

ETM (*n_f*=2, arXiv:1308.1851, JHEP 2014) vs. FNAL/MILC (*n_f*=3, arXiv:1602.03560)



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LQCD results for

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errors (in %) (preliminary) FLAG-3 averages + new results



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Nonfactorizable contributions:

$$\langle P\,\ell\ell|Q_i(y)|B\rangle \sim \bar{u}_\ell \gamma^\mu v_\ell \int d^4x \, e^{iq(x-y)} \langle P|TJ^\mu_{\rm em}(x)Q_i(y)|B\rangle$$

• at low recoil (high q^2) use OPE (Grinstein&Pirjol, hep-ph/0404250, PRD 2004; others): expand in $1/q^2 \sim 1/m_b^2$

 $\langle P \, \ell \ell | Q_i | B \rangle \sim f_{+,0,T}$ + quark-hadron duality violations (Beylich et al, arXiv:1101.5118, EPJC 2011)

• at high recoil (low q^2) use SCET: $E_P \sim m_b/2$, expand in Λ/M_B

$$\begin{array}{c} \langle P\,\ell\ell|Q_i|B\rangle \sim f_{+,0,T} + \ \phi_B \star T \star \phi_P \\ & \uparrow \\ \hline \\ \text{hard} \sim m_b^2 \end{array} \quad \begin{array}{c} \text{hard collinear} \sim \Lambda \ m_b^2 \end{array}$$



Experiment vs. Theory







Experiment vs. Theory







Experiment vs. theory

- LHCb data + FNAL/MILC form factors (arXiv:1509.00414, JHEP 2015;1403.8044, **JHEP 2014**)
- focus on large bins above and below charmonium resonances
- theory errors commensurate with experiment
- yields $\sim 1-2\sigma$ tensions
- \Rightarrow determine $|V_{td}/V_{ts}, |V_{td}|, |V_{ts}|$

or constrain Wilson coefficients



D. Du et al (arXiv:1510.02349, PRD 2016)





in the SM: $A_{\rm FB} = 0$ $F_H^\ell \sim m_\ell^2/M_B^2$

D. Du et al (arXiv:1510.02349, PRD 2016)





Phenomenology for $\Lambda_b \to \Lambda \ell^+ \ell^-$

Experiment vs. theory

- LHCb data + Detmold&Meinel form factors (arXiv:1503.07138, JHEP 2015)
- focus on regions above and below charmonium resonances
- exp. data lie above SM theory $\sim 1-3\sigma$ tensions

Detmold & Meinel (arXiv:1602.01399, PRD 2016)





theoretically clean



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B Mixing and FCNC decays



*from CKMfitter 2015 (hep-ph/0406186, http://ckmfitter.in2p3.fr)

Exclusive vs. inclusive $|V_{cb}|$ and $|V_{ub}|$



UT analysis



UT analysis



BSM phenomenology for $B \to K, \pi \ \ell^+ \ell^-$

Constraints on Wilson coefficients (C_9 , C_{10})

 \Leftrightarrow New physics contributions modify the Wilson coefficients:

 $C_i \to C_i + C_i^{\rm NP}$

at the high scale, $\mu_0 = 120 \text{ GeV}$

★ take $C_{7,8}^{\text{NP}} = 0$ using constraints from $B \to X_s \gamma$ ★ assume MFV so that $C_i(b \to s \,\ell \ell) = C_i(b \to d \,\ell \ell)$ ★ assume $C_{9,10}^{\text{NP}}$ are real (no new CP violating phases) ★ take measured $\Delta \mathcal{B}(B \to K, \pi \,\mu^+ \mu^-)$ in $\Delta q^2 = 1 - 6, 15 - 22 \text{ GeV}^2$ ★ and FNAL/MILC form factors

 \bigstar add $B_s \rightarrow \mu^+ \mu^-$ constraint with lattice f_{Bs}



Constraints on Wilson coefficients (C_9 , C_{10})



D. Du et al (arXiv:1510.02349, PRD 2016)

BSM phenomenology for $B \to K, \pi \ \ell^+ \ell^-$

Constraints on Wilson coefficients (C_9 , C_{10})



D. Du et al (arXiv:1510.02349, PRD 2016)

, NP,





Horgan et al (arXiv:1310.3887, PRL 2014; arXiv:1310.3722, PRD 2014, arXiv: 1501.00367)

Theoretical framework for weak decays to resonances, $B \to K \pi \, \ell \ell$ etc... being developed (Briceño et al, arXiv:1406.5965, PRD 2015; Agadjanov et al, arXiv:1605.03386)

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BSM phenomenology: LFU τ/ℓ



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SM prediction for
$$R(\pi) = \frac{\mathcal{B}(B \to \pi \tau \nu_{\tau})}{\mathcal{B}(B \to \pi \ell \nu)} = 0.641(17)$$

A. El-Khadra



BSM phenomenology: LFU μ/e

Lepton universality test: $B \to K \mu^+ \mu^- / B \to K e^+ e^-$







~2.6 σ tension between LHCb measurement and SM theory

In the SM these ratios are insensitive to the form factors (see also C. Bouchard et al, arXiv:1303.0434, PRL 2013)

Summary

ightarrow Precise LQCD results for semileptonic form factors for *B* → *π*, *K*, *D* transitions (and $\Lambda_b \rightarrow \Lambda, \Lambda_c, p$ transitions)

SM pre/postdictions with errors that are commensurate with experimental uncertainties

- $\blacksquare \sim 2\sigma$ tensions between exp. and SM theory
- \Rightarrow new LQCD results for neutral *B* meson mixing matrix elements with significantly smaller uncertainties
 - \rightarrow emerging ~2 σ tensions between loop processes and CKM unitarity
- \Rightarrow observed LFU violating effects ~2-4 σ level:
 - ♦ still need LQCD form factors for $B \rightarrow D^*$ at nonzero recoil
 - \bullet also needed for exclusive $|V_{cb}|$ determination

Thank you!

Farah Willenbrock

Backup slides

Introduction to Lattice QCD

 $\langle \mathcal{O} \rangle \sim \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A \mathcal{O}(\psi,\bar{\psi},A) e^{-S} \qquad \qquad S = \int d^4x \left[\bar{\psi}(\not\!\!D+m)\psi + \frac{1}{4} (F^a_{\mu\nu})^2 \right]$

use monte carlo methods (importance sampling) to evaluate the integral.

Note: Integrating over the fermion fields leaves det(D + m) in the integrand. The correlation functions, O, are then written in terms of $(D + m)^{-1}$ and gluon fields.

steps of a lattice QCD calculation:

- 1. generate gluon field configurations according to $det(D+m) e^{-S}$
- 2. calculate quark propagators, $(D+m_q)^{-1}$, for each valence quark flavor and source point
- 3. tie together quark propagators into hadronic correlation functions (usually 2 or 3pt functions)
- 4. statistical analysis to extract hadron masses, energies, hadronic matrix elements, from correlation functions
- 5. systematic error analysis

...of lattice spacing, chiral, and finite volume effects is based on EFT (Effective Field Theory) descriptions of QCD → ab initio

The EFT description:

- provides functional form for extrapolation (or interpolation)
- Can be used to build improved lattice actions/methods
- Solution of the size of systematic effects and the size of systematic effects are specified as the specified

discretization effects



discrete space-time \rightarrow discrete QCD action Symanzik EFT: $\langle \mathcal{O} \rangle^{\text{lat}} = \langle \mathcal{O} \rangle^{\text{cont}} + O(ap)^n$ *p* is the typical momentum scale associated with $\langle \mathcal{O} \rangle$ for light quark systems, $p \sim \Lambda_{\text{QCD}}$



a (fm)

The form of $O(ap)^n$ depends on the details of the lattice action.

All modern light-quark actions start at n = 2 (improved Wilson, twisted-mass Wilson, asqtad, HISQ, Domain Wall, Overlap, ...).



for charm lattice spacings are sufficiently small so that we can use improved light quark methods

- avoid errors of $(am_b)^2$ by using EFT in the formulation/matching of lattice action/currents:
 - relativistic HQ actions (Fermilab, Columbia, Tsukuba)
 - + HQET
 - + NRQCD

or

- use the same improved light quark action as for charm (HISQ, twisted mass Wilson, NP imp. Wilson, Overlap, ...)
 - + keep $am_h < 1$
 - + use HQET and/or static limit to extrapolate to the physical b quark mass

light quark mass effects

Simulations with $m_{\text{light}} = 1/2 (m_u + m_d)$ at the physical u/d quark masses are now available, but many results still have

 $m_{\rm light} > 1/2 (m_u + m_d)_{\rm phys}$

 χ PT can be used to extrapolate/interpolate to the physical point.

Solution Θ Can include discretization effects (for example, staggered χ PT)

It is now common practice to perform a combined continuum-chiral extrapolation/interpolation

finite volume effects

One stable hadron (meson) in initial/final state:

```
If L is large enough, FV error \sim e^{-m_{\pi} L}
```

 Θ keep $m_{\pi} L \gtrsim 4$

To quantify residual error:

- Solution compare results at several *L*s (with other parameters fixed)

The story changes completely with two or more hadrons in initial/final state! (or if there are two or more intermediate state hadrons)

other effects

- ✓ statistical errors: from monte carlo integration consider/include systematic errors from correlator fit procedure
- ✓ n_f dependence: realistic sea quark effects: use $n_f = 2+1$ or $n_f = 2+1+1$ Note: $n_f = 2$ (effects due to quenching the strange quark appear to be small)
- renormalization (and matching):
 - ⇒ with lattice perturbation theory: need to include PT errors
 - ⇒ nonperturbative methods
 - ⇒ use absolutely normalized currents where possible

...of lattice spacing, chiral, and finite volume effects is based on EFT (Effective Field Theory) descriptions of QCD → ab initio

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- provides functional form for extrapolation (or interpolation)
- Solution can be used to build improved lattice actions/methods
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To control and reliably estimate the systematic errors repeat the calculation on several lattice spacings, light quark masses, spatial volumes, ...

CKM determinations

 V_{ud} V_{us} V_{ub} $\pi \rightarrow \mu v$ $K \rightarrow \pi \ell \nu \quad B \rightarrow \pi \ell \nu, B_s \rightarrow K \ell \nu$ $K \rightarrow \mu \nu$ $\Lambda_b \rightarrow p \ell v$ V_{cd} V_{cb} V_{cs} $V_{cs} \qquad V_{cb}$ $D \rightarrow K \ell v \qquad B_{(s)} \rightarrow D_{(s)}, D^*_{(s)} \ell v$ $D \rightarrow \pi \ell v$ $D \rightarrow \ell \nu$ $D_s \rightarrow \ell v$ $\begin{array}{ccc}
 & & V_{ts} \\
B^0 - \overline{B^0} & B_s^0 - \overline{B_s^0} \\
B \rightarrow \pi \ell \ell & R \rightarrow T
\end{array}$ V_{tb} $(
ho,\eta)$ $oldsymbol{K}^0-\overline{oldsymbol{K}^0}$

The *z*-expansion



The form factor can be expanded as:

$$f(t) = \frac{1}{P(t)\phi(t,t_0)} \sum_{k=0}^{\infty} a_k(t_0) z(t,t_0)^k$$

Bourrely at al (Nucl.Phys. B189 (1981) 157) Boyd et al (hep-ph/9412324,PRL 95) Lellouch (arXiv:hep- ph/9509358, NPB 96) Boyd & Savage (hep-ph/9702300, PRD 97) Bourrely at al (arXiv:0807.2722, PRD 09)

- P(t) removes poles in $[t_{-},t_{+}]$
- The choice of outer function ϕ affects the unitarity bound on the a_k .
- In practice, only first few terms in expansion are needed.

B meson decay constant summary S. Aoki et al (FLAG-2 review, arXiv:1310.8555, FLAG-3 update)





CMS+LHCb combined (arXiv:1411.4413, Nature 2015)



SM predictions depend on $f_{B(s)}$ or \hat{B}_{B_s}



CMS+LHCb combined (arXiv:1411.4413, Nature 2015) and ATLAS (arXiv:1604.04263)



SM predictions depend on $f_{B(s)}$ or \hat{B}_{B_s}

D meson summary

errors (in %) (preliminary) FLAG-3 averages + new results $f_{D_{s}}/f_{D^{+}}$ small errors due to + physical light quark masses $(f_{D(s)})$ f_{D_s} improved charm-quark action f_{D^+} ensembles with small lattice spacings PCAC or NPR $f_+^{DK}(0)$ $f_{+}^{D\pi}(0)$ \hat{B}_D^i • First results for *D* mixing bag parameters (all five) of local operators by ETM (2013, 2014) $n_f = 2, 2+1+1$ 3 2 0 4 • work in progress: error in % FNAL/MILC (J. Chang thesis), see backup

slides